
Evaluation of One- and Two-Pole-Pair Slotless Bearingless Motors With Toroidal Windings

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Evaluation of One- and Two-Pole-Pair Slotless Bearingless Motors With Toroidal Windings

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Abstract—In this paper, winding topologies for one- and two-pole-pair rotors are analyzed and compared for a slotless bearingless disk drive with toroidal windings. The basis of the studies is a six-phase motor with a diametrically magnetized one-pole-pair rotor. Due to the absence of mechanical bearings and the significantly large air-gap capability, the motor is suitable for applications with high purity and special chemical demands. Its slotless design results in low losses even at high rotational speeds. To improve the operational behavior of the rotor in different applications, the influence of higher pole pair numbers on the passive bearing stiffness is examined. A possible winding configuration for these rotors is proposed and evaluated for their bearing and motor performance. Based on the results, a further prototype was built and is presented in this paper.

Index Terms—Active magnetic bearing, bearingless motor, high-speed drive, slotless motor.

I. INTRODUCTION

In a bearingless permanent-magnet (PM) synchronous motor, the rotor is spatially suspended and rotated without any mechanical contact [1]. Due to the lack of lubrication and abrasion, as well as the possibility to hermetically isolate the rotor and stator from the environment, bearingless motors and magnetically suspended motors are used, for example, in semiconductor manufacturing [2], in pumps or mixers for delicate fluids [3]–[7], or in high-speed applications [8], [9].

Due to the absence of mechanical contact between the rotor and stator, mechanical bearings and shaft feedthroughs can be omitted. This allows for a hermetic encapsulation of the rotor and stator, which results in high resistance against chemically aggressive fluids and gases and in increased lifetime compared with conventional motors. The absence of abrasion and lubrication enables operation in high-purity applications.

In the presented slotless bearingless motor concepts, a ring-shaped rotor is magnetically levitated without mechanical contact in the middle of an annular stator. This disk topology results in enhanced compactness due to its low axial length. The coils, which are toroidally wound on the stator iron, can generate both bearing forces and drive torque. This slotless motor topology has the advantage of reduced losses at high rotational and circumferential speeds compared with conventional slotted machines.

Especially when big air gaps are required, for example, due to a chemical sealing, bearingless motors have significantly lower bearing stiffness than mechanical ones, which limits possible applications. In pumps and blowers, for instance, axial and radial forces resulting from the differential pressure of the fluid act on the impeller. This means a considerable movement of the magnetically levitated rotor and limits the achievable pump, blow, or turbine performance [10].

Motors with a one-pole-pair diametrical magnetized rotor and six or ten coils are already proposed in [11]–[13]. To improve both the passive bearing stiffness and the active bearing and motor performance of the bearingless motor, different rotor magnetizations, rotor pole pair numbers, and winding configurations are necessary. Fig. 1 shows the two topologies that are evaluated. The coils are enumerated to show the connection

Fig. 1. Principle setup of two different slotless bearingless disk drives. Novel topology B features a two-pole-pair rotor and therefore has improved passive bearing stiffness. Topology A has combined coils and generates force and torque simultaneously with the same set of coils as a sum of motor and suspension currents will be supplied. (a) Topology A. (b) Topology B.
of the coils in the schematics, which will be shown later. Topology A has six coils and a one-pole-pair diametrical rotor, as presented by the authors in [11]. This is the best topology when low losses are demanded. Novel topology B uses a two-pole-pair rotor and therefore features higher passive stiffness values, which can be important, for example, in blower or pumping applications. The two topologies will be analyzed and compared in this paper.

It has to be distinguished between combined coils, where drive and bearing currents are superposed on one set of coils, and separate coils for motor and bearing. In [14] and [15], a slotted bearingless motor with combined windings is presented, and it is shown that combined windings lead to a more compact and simple winding scheme. However, with 3 three-phase inverters, the power electronics are quite extensive.

The goal is now to find a possible winding scheme with low complexity in winding and power electronics for rotors with two-pole pairs. We will show that, with a two-pole-pair rotor, the passive bearing stiffness is significantly increased.

At first, the magnetic air-gap field distributions of rotors with different pole pair numbers are evaluated. The influence of the pole pair number on the passive bearing stiffness values is shown.

As only specific winding configurations are capable of generating torque and force independently for given rotor magnetization, we derive the criteria for developing suitable configurations for separated and combined winding schemes. The winding concepts differ with regard to the coil number and the interconnection of the coils.

By 3-D finite-element (FE) simulations, the selected winding topologies are compared regarding their bearing and motor capabilities, as well as passive stiffness values. Additionally, possible loss mechanisms and the influence of the pole pair number onto the losses are discussed. At the end, a prototype of topology B with a two-pole-pair rotor will be presented and compared with the prototype of topology A. It will be shown that novel topology B is preferable in applications, where high bearing stiffness values are necessary. However, lower losses at high speeds can be achieved with topology A.

II. WORKING PRINCIPLE OF THE SLOTLESS BEARINGLESS DISK DRIVE

In [11], the authors present the detailed working principles and simulation and test results of the bearingless slotless disk drive with six coils and a diametrically magnetized one-pole-pair rotor. The presented prototype is the basis for the topology evaluation in this paper.

A. Passive Bearing Stiffness Values

The disk-type bearingless motor is passively stable in the tilting and axial directions and has to be actively controlled in the radial direction. Fig. 2 shows the magnetic field that is generated by the PMs on the rotor for a one-pole-pair (P1) and a two-pole-pair (P2) magnetization. For the P1 magnetization, the axes of magnetization \((d\text{-axis})\) and the orthogonal direction \((q\text{-axis})\) are highlighted. This magnetic field generates reluctance forces between the stator iron and the rotor. Therefore, deflection in the axial \(z\)-direction will lead to a counteracting force

\[
dF_z = -c_z \cdot dz
\]

due to the axial stiffness constant \(c_z\) similar to a spring constant. With this definition, a positive value of the stiffness leads to a stabilizing force.

Similarly, tilting of the rotor results in a counteracting torque. At a rotor with P1 magnetization, it has to be distinguished between the rotation around the axis of magnetization (angle \(\alpha\)) and the rotation perpendicular to it (angle \(\beta\)). The tilting stiffness values \(c_\alpha\) and \(c_\beta\) are then defined as

\[
c_\alpha = -\frac{dT_\alpha}{d\alpha} \quad \text{and} \quad c_\beta = -\frac{dT_\beta}{d\beta}.
\]

When the rotor is displaced from its center position in the radial direction, a radial force will act in the same direction. Therefore, the radial stiffness values

\[
c_d = -\frac{dF_x}{dx_d} \quad \text{and} \quad c_q = -\frac{dF_y}{dx_q}
\]

are destabilizing. Here, \(x_d\) means a displacement in the direction of the magnetization and \(x_q\) means a displacement perpendicular to the magnetization.

A one-pole-pair rotor results in anisotropic radial and tilting stiffness. Deflection of the rotor in the direction of magnetization \((d\text{-axis})\) results in a higher attractive force than deflection perpendicular to the magnetization \((q\text{-axis})\). The same holds for the tilting stiffness. Rotation around the \(d\)-axis results in a lower torque than rotation around the \(q\)-axis. This results in a broad resonance frequency range, as shown in [16]. Rotors with higher pole pair numbers, however, show almost isotropic stiffness values in all directions

\[
P = 1 : c_\alpha < c_\beta \quad \text{and} \quad c_q < c_d
\]

\[
P \geq 2 : c_\alpha \approx c_\beta \quad \text{and} \quad c_q \approx c_d.
\]

Therefore, rotor magnetizations of two or more pole pairs are favorable for increased bearing stability.
Fig. 3. Simulated armature reaction field due to the coil currents indicated by different shades of blue and red. Depending on the type of the current, each topology generates the corresponding field for force or torque generation. The rotor is not magnetized in this simulation to only show the armature field. (a) Topology A-P1 with drive current. (b) Topology A-P1 with bearing current. (c) Topology B with drive current. (d) Topology B with bearing current.

B. Bearing Force and Drive Torque Generation

The stator of the bearingless motor has to produce both suspension force and drive torque simultaneously. For a resulting torque acting on the rotor, the stator coils have to generate a magnetic armature reaction field in the magnetic gap with a pole pair number

$$p_{drv} = p$$  \hspace{1cm} (5)

which is equal to the pole pair number $p$ of the rotor magnetic field. A force will be generated by a stator field with a pole pair number

$$p_{bng} = p \pm 1$$  \hspace{1cm} (6)

as described in [17].

Fig. 2 shows the simulated flux plots of rotors P1 and P2 with no currents in the coils. Fig. 3 shows the simulated armature reaction field of topologies A and B. The rotor is not magnetized in this simulation to only show the stator field. It shows that, depending on the current scheme, topology A produces a one-pole-pair field with a drive current [see Fig. 3(a)] and a two-pole-pair field with a bearing current [see Fig. 3(b)]. This correlates to the one-pole-pair rotor field. Respectively, topology B generates a $p_{drv} = 2$ field for the drive and a $p_{bng} = 3$ field for the bearing corresponding to the two-pole-pair rotor.

In Fig. 4, the superposition of the PM rotor field and armature reaction field resulting from a bearing current is exemplarily shown for topology A. This corresponds to a combination of Fig. 2(a) with Fig. 3(b). In this case, the armature reaction field is superposed with the rotor field so that the field density is increased at the right side of the rotor and decreased at the left side. This leads to a force on the rotor toward the stronger field density on the right. Additionally, Lorentz forces are generated by the rotor magnetic field that penetrates the coils in the air gap. The current in the coils generates a force in the tangential direction. Summing up all Lorentz forces in the coils leads to a bearing force in the same direction as the reluctance force. Together, Lorentz and reluctance forces add up to the total bearing force. As shown in [11], the contribution of both force components is approximately equal for the case at hand.

However, this is only true if $p_{bng} = p + 1$. For an armature reaction field with a pole pair number smaller than that of the rotor ($p_{bng} = p - 1$), both forces point to opposite directions and partially cancel each other out, as proven in [2]. This would lead to reduced bearing forces, as will be also shown in the 3-D simulation results in Section IV. Therefore, to achieve high bearing forces, a magnetic air-gap field with $p_{bng} = p + 1$ has to be generated by the set of coils.

By analytic calculation, it can be shown that the bearing force $F_{bng}$ and the drive torque $T$

$$F_{bng} = k_F \cdot \hat{I}_B \left( \frac{\cos(\varphi_F)}{\sin(\varphi_F)} e_x + e_y \right)$$  \hspace{1cm} (7)

$$T = k_T \cdot \hat{I}_D$$  \hspace{1cm} (8)

are proportional to the amplitude of bearing current $\hat{I}_B$ or drive current $\hat{I}_D$ [11]. The force coefficient $k_F$ and the drive coefficient $k_T$ can be determined by a finite-element method (FEM) simulation. The force direction can be controlled by the phase shift $\varphi_F$ between electrical rotor angle and bearing current.

III. EVALUATION OF PASSIVE BEARING PROPERTIES FOR DIFFERENT MAGNET CONFIGURATIONS

Depending on the application, considerably large magnetic gaps are necessary, for instance, if a chemically resistant wall...
has to be inserted between the rotor and stator. This drastically reduces the passive stiffness values of the bearing, as shown in [12]. Therefore, a strong magnetic field in the air gap between magnet and stator iron is important. Moreover, the active force and torque generation depends particularly on the fundamental wave of the flux density distribution in the air gap.

To compare the different magnetizations, a fixed geometry, which was already used in [11], is chosen. The dimensions of stator and magnet, which are the most important dimensions for the evaluation, are shown in Table II (prototype A). This geometry results from an optimization with a one-pole-pair rotor and yields a significantly large air gap, suitable for pump and blower applications.

Fig. 5 shows a diametrical one-pole-pair (P1) and a two-pole-pair (P2) magnet configuration, as well as the harmonic analysis of the magnetic field distribution in the middle of the magnetic gap, which was simulated with a FEM simulation. The diagram below shows that the diametrical magnetization yields a purely sinusoidal magnetic air-gap field. This is the best for a good bearing and drive performance, as well as for lowest losses, as higher harmonics do not contribute to force and torque but generate losses. However, as previously shown, the stiffness values of rotor P1 are anisotropic, which makes it necessary to evaluate higher pole pairs.

For higher pole pair numbers, the magnetization [see Fig. 5(b)] consists of alternately inward and outward magnetized magnets. It yields a higher flux density but also higher harmonics in the field distribution.

The axial stiffness values, which are simulated with a 3-D magnetostatic FEM simulation tool, are shown for different pole pair numbers and magnet configurations in Fig. 6(a). From rotor P1 to rotor P2, the axial stiffness increases by 39%.

With pole pair numbers higher than \( p = 2 \), the stiffness would decrease again.

Similar results are found for the tilting stiffness, which are shown in Fig. 6(b). The tilting stiffness increases by 37% with a two-pole-pair magnetization (P2) compared with the mean tilting stiffness of the one-pole-pair rotor (P1). In Fig. 6(b), the minimum and maximum tilting stiffness are also shown for rotor P1 that results from the aforementioned anisotropy. As the tilting stiffness values are isotropic for the \( p \geq 2 \) magnetizations, rotor P2 has an increased tilting stiffness value of 300% compared with the minimal value of rotor P1.

Summarizing, the highest passive bearing stiffness values can be achieved with a two-pole-pair rotor. A further increase in the pole number reduces the stiffness values.

**IV. WINDING CONCEPTS**

For each pole pair number, it is necessary to derive a winding configuration that is capable of generating both torque and force independently.

**A. Winding Criteria for Combined Coils**

With combined coils, both torque and force are generated by the same set of coils. Because of the different pole numbers for bearing and drive, the winding scheme contains no repetitive elements. This means that the number of phases \( m \) is equal to the number of coils \( N \). However, one coil can be separated in two coils with reversed winding direction, which are connected in series, as proposed in [16]. This will have no direct effect on the feasibility of the winding configuration and is not examined in this paper.

To generate a field with a given pole pair number \( p \), the number of coils

\[
N \geq 2 \cdot p
\]

has to be at least twice as high to avoid aliasing. For the bearing, this is even stricter as there are two degrees of freedom that have to be actively controlled. If the number of poles and coils would be equal, the bearing phases would have a phase shift of
180° and therefore only one degree of freedom. Altogether, the number of coils for the bearing

\[ N > 2 \cdot p_{bng} \]  

has to be more than twice as high as the number of pole pairs needed for bearing operation.

For a one-pole-pair rotor, a bearing field with \( p = 2 \) has to be generated. This leads to a minimal coil number of five. As aforementioned, five- and six-phase motors with a one-pole-pair rotor already exist [11]–[13]. The six-coil topology is referred to as “topology A” in this paper.

Theoretically, this topology is also possible with a two-pole-pair rotor as it can generate a \( p_{bng} = 2 \) stator field for the drive and a \( p_{bng} = 1 \) (\( p_{bng} = p - 1 \)) field for the bearing. As aforementioned, this will lead to a low bearing performance. This topology is referred to as “topology A-P2” and will be tested for comparing the losses of one- and two-pole-pair machines, but due to the poor bearing performance, it is not recommended for any application.

Therefore, with a two-pole-pair rotor, a bearing field with \( p_{bng} = 3 \) has to be generated by the coils. A six-coil motor is not able to generate a continuous \( p_{bng} = 3 \) stator field as the phases would have a 180° phase shift. With the criteria in (10), at least seven coils are needed for a two-pole-pair rotor. A topology with eight combined coils, which results in less complicated power electronics than with seven coils, is presented by the authors in [18]. As this topology is still quite complicated regarding power electronics and wiring, it will not be discussed here.

B. Winding Criteria for Separated Coils

With separated coils, the magnetic field for each bearing and motor operation is generated by an independent set of coils. Therefore, two independent inverters can be used. Then, the winding configuration consists of repetitive elements corresponding to the respective pole pair number of bearing and motor field.

For the bearing field, at least two phases per pole pair are necessary, as two degrees of freedom have to be controlled. To maintain manufacturability, the coil numbers for bearing and motor should be the same. Therefore, winding configurations with 9, 12, or 18 coils each for bearing and motor are possible.

Simulations of the \( 2 \times 9 \)-coil topology showed that, due to adverse field harmonics of the armature reaction field, a coupling of the bearing and the drive system exists. Therefore, this topology will not be explained here, but it is shown in [18]. The \( 2 \times 12 \)-coil topology has three phases for the drive and two phases for the bearing. This results in a higher power electronic effort than a solely three-phase system.

The \( 2 \times 18 \)-coil topology does not show this disadvantageous behavior and has a three-phase system each for drive and bearing. It is referred to as “topology B” and will be further examined in this paper.

C. Winding Connection

In Fig. 7, the winding connection of the aforementioned motor topologies is shown. The numbering of the coils is shown in Fig. 1. Every coil is wound in the same direction. A negative sign in the winding scheme indicates that the coil connection is reversed. For each topology, we define two independent sets of currents. With the drive currents \( I_{D,U}, I_{D,V}, \) and \( I_{D,W} \), the drive torque is generated, whereas the bearing currents \( I_{B,U}, I_{B,V}, \) and \( I_{B,W} \) generate the bearing force. Each current system consists of three 120° phase-shifted currents.

As in combined windings the coils generate both force and torque, the drive and bearing currents have to be superposed on each phase. For topology A-P2, bearing and drive current have to be swapped. With the separate coils of topology B, there is one designated star system each for bearing and drive. Every coil is wound in the same direction. A negative sign in the winding scheme indicates that the coil connection is reversed. In Fig. 7, the connection of the coils to the power electronics is shown for both topologies. The winding connection is exemplarily shown for topology A.

For the separate winding concept, the drive currents are directly fed into the three-phase system of the drive coils and similarly with the bearing currents.

The connection of the coils to the power electronics is shown in Fig. 8. The power electronics consist of two inverter modules with six switches each to power two three-phase systems.
D. Simulation Results of the Proposed Topologies

Subsequently, the two topologies have been simulated in 3-D FE simulations. The results are shown in Table I. For comparison reasons, the bearing and drive constants are calculated, assuming a fill factor of 1 in the coils.

It shows that the two-pole-pair rotor has significantly higher passive stiffness values. As expected, the active bearing force of topology A-P2 is significantly lower. Together with the higher destabilizing radial stiffness, the magnetic bearing performance of topology A-P2 is not recommended for real applications.

However, with topology B, the highest bearing force is achievable, which is sufficient to compensate also the high radial stiffness of the two-pole-pair rotor.

Therefore, with novel two-pole-pair topology B, a promising alternative to one-pole-pair topology A-P1 is developed. The passive stiffness values and the active bearing and drive performances exceed the values of topology A-P1.

V. Test Results

A. Prototypes

For all proposed topologies, two prototypes were built and tested. Prototype A refers to topology A and can be operated with a P1 and a P2 rotor. This allows direct comparison of stiffness and losses for both rotor types. Since prototype A-P2 shows, as expected, very week bearing performance, a further prototype B with topology B was built.

Prototype “A-P1” with a diametrical magnetized one-pole-pair rotor and six combined coils has been already built and successfully tested [11]. Fig. 9 shows the prototype mounted on a plate together with a printed circuit board for sensor signal amplification. Stable operation is possible for up to 20 000 r/min, which is the mechanical limit of the rotor.

When operating this motor with a two-pole-pair rotor (prototype A-P2), the startup bearing current needed to start the levitation is around 2.7 times higher than with the P1 rotor. During operation, the rotor has to stay perfectly in the radial center so that the bearing does not have to generate high forces to compensate for the position offset. Then, stable operation was possible for up to 17 000 r/min.

Additionally, a second prototype (“prototype B”) with a two-pole-pair rotor and 2 × 18 separated coils corresponding to topology B is built and depicted in Fig. 10. As higher losses occur with two-pole-pair rotors at the same rotational speed, prototype B was built with a significantly bigger rotor outer diameter so that higher circumferential speeds can be achieved with lower rotational speeds.

While prototype A has a rotor diameter of 102 mm, prototype B has a rotor size of 164.8 mm. The bigger size also enhances the manufacturability, particularly with the high number of coils that have to be placed on the stator of prototype B. The geometrical details of both prototypes are compared in Table II.
With prototype B, rotational speeds of up to 12000 r/min are reached. This is the limit of mechanically safe operation. The prototype showed stable bearing behavior in standstill and during rotation. Rigid body resonances were observed only between 1500 and 2000 r/min, whereas prototype A-P1 is instable in the range of 500–2000 r/min. This can be explained with the anisotropic radial and tilting stiffness values of prototype A-P1. When the rotor is radially displaced or tilted, the attractive force or torque toward the stator will vary during rotation if the stiffness is not isotropic. This can easily trigger oscillations. Additionally, the simulated force, which was used to calculate the stiffness values in Table I, is shown for each prototype by the applied weight, in dependence on the rotor displacement. The gradient of the axial force correlates to the stiffness of the bearing coefficient, the rotor was levitated at standstill with a lever that pulls on a force measurement device, whereas the drive current was measured. The radial stiffness can be measured by displacing the rotor and measuring the bearing current. As the bearing performance of prototype A-P2 is too bad to hold the rotor weight in horizontal position, these measurements were not possible with this motor. The results are shown in Table III. The values in brackets show the simulation results.

The axial stiffness measurement of two-pole-pair rotor P2 showed an increase of 50% compared with rotor P1 in prototype A. Due to the increased size, the stiffness of prototype B is even 43% higher than that of prototype A-P2.

The simulation results fit very well to the measurements of stiffness and force and drive constant. Therefore, not only the impact of the pole pair number on the stiffness is proven but also the correctness of the simulation used in the previous sections is confirmed.

### C. Losses

Increasing the pole pair number significantly affects the losses of the machine as the frequencies of magnetic fields and currents are increased. In particular, the main losses of this type of bearingless machine, which are losses in the stator iron and eddy-current losses in the copper, are strongly frequency dependent, as demonstrated in [19] and [20].

When the pole pair number is increased from one to two, there are different effects influencing the losses. As most loss components quadratically depend on the frequency, these losses will be increased by a factor of four. Additionally, the field distribution of a two-pole-pair rotor might show higher spatial harmonics, which was already shown in Fig. 5. This increases these losses even more.

Additionally, there exist losses that do not depend on the magnetic field, for example, controller losses, air friction losses at the rotor surface [21], losses due to carrier harmonics [22]–[24], or resistive copper losses depending on the load torque.

The losses in the machines were measured with a digital oscilloscope by multiplying instantaneous current and voltage and subsequent averaging over time. In Fig. 12, the measured losses of the prototypes are plotted against the rotational speed. As the maximum copper losses are below 2 W in both motors, they are not explicitly illustrated.

It shows that the losses of prototype A-P2 are higher than those of prototype A-P1 by a factor of around 2.2 due to the increased pole pair number. Due to the increase in size, the losses of prototype B are additionally increased by a factor of 2.1. At the same circumferential speed, prototype A-P1 produces the lowest losses.

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**Fig. 11.** Axial displacement versus axial force (symbols) measured and (lines) simulated for prototypes A and B.

**Table III**

<table>
<thead>
<tr>
<th>Prototype</th>
<th>A-P1</th>
<th>A-P2</th>
<th>B</th>
</tr>
</thead>
<tbody>
<tr>
<td>axial stiffness</td>
<td>8.1 (8.7)</td>
<td>12.2 (12.1)</td>
<td>17.4 (16.2)</td>
</tr>
<tr>
<td>radial stiffness</td>
<td>-16.6 (-18)</td>
<td>---</td>
<td>-27.6 (-27.2)</td>
</tr>
<tr>
<td>torque coefficient</td>
<td>0.144 (0.147)</td>
<td>---</td>
<td>0.188 (0.190)</td>
</tr>
<tr>
<td>force coefficient</td>
<td>2.97 (2.72)</td>
<td>---</td>
<td>4.34 (4.83)</td>
</tr>
</tbody>
</table>
Therefore, for high rotational speeds and high circumferential speeds, which is important for applications like pumps and blowers, a low pole pair number is preferable. In particular, the lower spatial harmonics in the rotor magnetization is a big advantage of the one-pole-pair rotor. However, there is still the possibility of reducing the losses of the two-pole-pair motor by improved magnetization and by using stranded litz wires and therefore omitting eddy current losses in the coils. This is will be an issue of further research.

VI. CONCLUSION

Summarizing, it can be said that, at disk-type motors, a two-pole-pair rotor will lead to significantly higher and isotropic passive bearing stiffness values. A novel coil topology has been proposed, which is possible with a two-pole-pair rotor. A motor with 36 coils was built to demonstrate the feasibility of this two-pole-pair machine. It showed a very stable operational behavior with a small resonance range and high isotropic stiffness values.

However, it also showed that the losses of a two-pole-pair rotor are roughly twice as high as with a one-pole-pair rotor. Especially when high speeds are required, the one-pole-pair rotor is better due to its low losses. If the bearing stiffness is not sufficient for the required application, which can be the case in pumps and blowers, a two-pole-pair motor can be chosen.

REFERENCES

Thomas Nussbaumer (S’02–M’06) was born in Vienna, Austria, in 1975. He received the M.Sc. (Hons.) degree in electrical engineering from Vienna University of Technology, Vienna, in 2001 and the Ph.D. degree from the Swiss Federal Institute of Technology Zurich (ETH Zurich), Zurich, Switzerland, in 2004.

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Prof. Kolar was a recipient of 21 IEEE TRANSACTIONS and conference prize paper awards, the 2014 SEMIKRON Innovation Award, the 2014 IEEE Power Electronics Society R. David Middlebrook Award, and the ETH Zurich Golden Owl Award for Excellence in Teaching.